18.354J / 3541

Nonlinear dynamics: continuum systems

Course intro



18.354J/3541 Nonlinear dynamics: continuum systems Spring 2020

Lectures:	MW 2:30-4.00	in	253 - 7826	
Instructor:	Jörn Dunkel			
Contact:	dunkel@mit.edu		253-7826 (office p	ohone)
Office Hours:	Thursdays 2:30-3:30 (2-381)			
Course website:	${\tt math.mit.edu}/{\sim}{\tt dunkel}/{\tt Teach}/18.354/$			

https://math.mit.edu/~dunkel/Teach/18.354/

course info, psets, lecture notes,

GRADING

- \bullet 60%: Problem sets
- 10%: Project proposal
- \bullet 30%: Final project presentation + report

HOMEWORK - PROBLEM SETS

Homework will be assigned roughly every two-three weeks. We aim to have 4 problem sets in total. Each homework set may contain analytical and computational problems, and even the odd experiment. Assignments must be handed in at the start of class on the due date.

First unexcused late homework score will be multiplied by 3/4. No subsequent unexcused late homework is accepted.

You are welcome to discuss solution strategies and even solutions, but please write up the solution on your own. Be sure to support your answer by explaining important steps or listing any relevant theorems. Be as clear and concise as possible. I strongly encourage the computational problems to be written in MATHEMATICA, MATLAB or JULIA.

TEXTBOOKS

Although there are no textbooks which follow the precise spirit of this course, there are literally hundreds of textbooks where the topics we will cover are discussed. For most lectures, typed notes can be downloaded from the course webpage. Useful books for further reading include:

- M. Cross, *Pattern Formation and Dynamics in Non-equilibrium Systems*, Cambridge University Press (2009).
- A. Goriely, The Mathematics and Mechanics of Biological Growth, Springer (2017).
- B. Audoly & Y. Pomeau, *Elasticity & Geometry*, Oxford University Press (2010).
- D. J. Acheson, *Elementary Fluid Dynamics*, Oxford University Press (1990).

IMPORTANT DATES

- \bullet Wed Feb 26 Problem Set 1 DUE
- \bullet Wed Mar 11 Problem Set 2 DUE
- Wed Mar 18 Proposal (1 page) for final project DUE
- Wed Apr 1 Problem Set 3 DUE
- Wed Apr 22 Problem Set 4 DUE
- May 3, 5 & 10 Final project presentations
- May 10 Final project report DUE

Note: The exact due dates for the P-sets may be subject to change

Course topic:

Physical Applied Mathematics

Q: What is Physical Applied Maths?



PAM is like cooking...



Often the ingredients (physical principles) are already known but not the way (mathematics/equations/couplings) to turn them into a nice dinner

With some creativity, many new dishes (novel phenomena) can be created (discovered/understood)

Why study Applied Maths?

- intellectual challenge
- obtain general understanding of physical phenomena and the world around us
- be able to make prediction about physical processes
- development of general tools to be applied to other fields

Discrete active matter





Bacterial biofilms



Drescher lab, MPI Marburg

PNAS 2016 Nature Physics 2019 Nature Microbiology 2019 Physical Review Letters 2019



Bacterial swarming





Drescher lab, MPI Marburg

PNAS 2019

Nature Physics 2018

Continuous biological (&) soft matter



Continuous active matter

Active fluids

Energy transport in active fluids



Geometric control of microbial suspensions



Goldstein lab, Cambridge

Cell membranes



Fakhri lab, MIT Physics

PNAS 2012 Physical Review Letter 2013 PNAS 2017 J Fluid Mech 2018 Physical Review Letters 2013 PNAS 2013 eLife 2014 Nature Physics 2016

Nature Physics 2020 (in press)

Historical backdrop ... quiz:

Some famous thought on Applied Math ?

'Eureka, Eureka'





c. 287 BC - c. 212 BC

'Mathematic is written for mathematicians'





19 February 1473 – 24 May 1543

'It would be better for the true physics if there were no mathematicians on earth'

8 February 1700 – 17 March 1782



<text><text><section-header><text><text><text><text><text><text><text><text>

'Now I will have less distraction'





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15 April 1707 – 18 September 1783
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"If this subjet has hitherto been considered from the wrong viewpoint and thus enveloped in mystery and surrounded by darkness, it is largely an unsuitable terminology which should be blamed. Had +1, -1 and $\sqrt{-1}$, instead of being called positive, negative and imaginary (or worse still, impossible) unity, been given the names say, of direct, inverse and lateral unity, there would hardly have been any scope for such obscurity."

'I do not know'



25 January 1736 - 10 April 1813



'Nature laughs at the difficulties of integration'



'Mathematicians are born, not made'





29 April 1854 – 17 July 1912

'Prediction is very difficult, especially about the future'



7 October 1885 – 18 November 1962



'It is more important to have beauty in one's equations than to have them fit experiment' and 'This result is too beautiful to be false'

8 August 1902 – 20 October 1984

'To those who do not know mathematics it is difficult to get across a real feeling as to the beauty, the deepest beauty, of nature'



May 11, 1918 – February 15, 1988

BSc 1939

Course content overview

Concepts

- linear/nonlinear PDEs for scalar, vector and tensor fields
- model derivation/construction, dimensional analysis
- solution techniques, Fourier methods
- stability, perturbation theory
- calculus of variations

Examples

- Hamiltonian dynamics & Liouville equation
- Brownian motion (free and confined)
- quantum mechanics (linear and nonlinear)
- pattern formation theory (Turing)
- elasticity theory
- fluid dynamics (passive and active)

Dimensional analysis



G I Taylor 1886-1975



Trinity nuclear test, July 1945 Life Magazine, August 20, 1945



Hamiltonian dynamics & Liouville equation

$$H = \sum_{i} \frac{p_i^2}{2m_i} + U(x_1, \dots, x_N)$$



Brownian motion



Random walks & diffusion



 $\frac{\partial \rho}{\partial t} = \alpha \frac{\partial^2 \rho}{\partial x^2}$



Mark Haw

David Walker

Quantum mechanics

$$i\psi_t = -\psi_{xx} + \omega^2 x^2 \psi$$





Higher-order quantum hydrodynamics

NLS GPE

$$i\hbar\partial_t \Psi(t,\mathbf{r}) = \left(-\frac{\hbar^2}{2m}\nabla^2 + g_0 |\Psi(t,\mathbf{r})|^2\right)\Psi(t,\mathbf{r})$$

$$u = g_0\delta(\mathbf{x} - \mathbf{x}')$$

 $\langle u_{\text{int}} \rangle = \frac{1}{2} \int \int d\mathbf{r}_1 \, d\mathbf{r}_2 \left[\Psi^*(t, \mathbf{r}_2) \Psi^*(t, \mathbf{r}_1) u_{\text{int}}(|\mathbf{r}_2 - \mathbf{r}_1|) \Psi(t, \mathbf{r}_1) \Psi(t, \mathbf{r}_2) \right]$

$$\hat{u}_{\text{int}}(k) = \sum_{j=0}^{\infty} g_{2j} k^{2j}$$
$$i\hbar \partial_t \Psi = \left[-\frac{\hbar^2}{2m} \nabla^2 + \left(\sum_{j=0}^{\infty} (-1)^j g_{2j} (\nabla^2)^j |\Psi|^2 \right) \right] \Psi$$

Analytically tractable theory of super-solids



Heinonen & JD, Physical Review A 2019

Pattern formation









zebra vs. granular medium vs. fingerprint



Mathematical theory of pattern formation

Philosophical Transactions of the Royal Society of London. Series B, Biological Sciences, Vol. 237, No. 641. (Aug. 14, 1952), pp. 37-72.

CHEMICAL BASIS OF MORPHOGENESIS THE

By A. M. TURING, F.R.S. University of Manchester

(Received 9 November 1951—Revised 15 March 1952)



 $\partial_t \boldsymbol{c}(t, \boldsymbol{x}) = \boldsymbol{R}(\boldsymbol{c}) + \boldsymbol{D} \cdot \boldsymbol{c}$

local (nonlinear)

non-local reactions diffusive interactions (linear)

Swift-Hohenberg theory

$$\partial_t \psi = -U'(\psi) + \gamma_0 \nabla^2 \psi - \gamma_2 (\nabla^2)^2 \psi$$

$$U(\psi) = \frac{a}{2}\psi^2 + \frac{b}{3}\psi^3 + \frac{c}{4}\psi^4$$





Beyond Turing

Starfish egg cell





Fakhri lab MIT Physics

Topological turbulence in the membrane of a living cell Tan*, Liu*, Miller*, Tekant, JD & Fakhri, Nature Physics (in press)

Calculus of variations

$$\begin{split} \frac{\delta I[Y]}{\delta Y} &= \lim_{\epsilon \to 0} \frac{1}{\epsilon} \left\{ I[f(x) + \epsilon \delta(x - y)] - I[f(x)] \right\} \\ &= \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial Y} \delta(x - y) + \frac{\partial f}{\partial Y'} \delta'(x - y) \right] dx \\ &= \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial Y} - \frac{d}{dx} \frac{\partial f}{\partial Y'} \right] \delta(x - y) dx. \end{split}$$

$$0 = \frac{\partial f}{\partial Y} - \frac{d}{dx} \frac{\partial f}{\partial Y'}$$

Elasticity



Elasticity



photo: Andrej Kosmrlj

Surface tension





Goldstein lab, Cambridge

Large drop in microgravity



Hydrodynamics

$$\int_{V} \frac{\partial \rho}{\partial t} dV = -\int_{S} \rho \mathbf{u} \cdot \mathbf{n} dS = -\int_{V} \nabla \cdot (\rho \mathbf{u}) dV. \qquad \qquad \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0.$$

$$\int_{V(t)} \rho \frac{D\mathbf{u}}{Dt} dV = \int_{V(t)} (-\nabla p + \rho \mathbf{g}) dV \qquad \qquad \frac{D\mathbf{u}}{Dt} = \frac{-\nabla p}{\rho} + \mathbf{g}.$$

Typical Reynolds numbers



~ meters



What happens at low Reynolds numbers ?



Low-Re (laminar) flow



Swimming at low Reynolds number

Navier - Stokes : $\langle \mathcal{R} \sim UL\rho/\eta \ll 1$ 1f

Time doesn't matter. The pattern of motion is the same, whether slow or fast, whether forward or backward in time.

The Scallop Theorem





Geoffrey Ingram Taylor





James Lighthill

Edward Purcell

 $0 = \mu \nabla^2 \boldsymbol{u} - \nabla p + \boldsymbol{f},$ $0 = \nabla \cdot \boldsymbol{u}.$ + time-dependent BCs

Superposition of singularities

stokeslet



$$p(\mathbf{r}) = \frac{\hat{\mathbf{r}} \cdot \mathbf{F}}{4\pi r^2} + p_0$$
$$v_i(\mathbf{r}) = \frac{(8\pi\mu)^{-1}}{r} [\delta_{ij} + \hat{r}_i \hat{r}_j] F_j$$

2x stokeslet = symmetric dipole



rotlet



 r^{-2}

flow ~ r^{-1}



E coli (non-tumbling HCB 437)





Drescher et al (2011) PNAS

E coli (non-tumbling HCB 437)





Drescher et al (2011) PNAS

Nonlinear HD: Solitons

KdV equation

$$\partial_t \phi + \partial_x^2 \phi + 6\phi \partial_x \phi = 0$$



Solitons



credit: Christophe Finot

Active fluids

Collective non-equilibrium dynamics in a multicellular system

Bacillus subtilis



Vortex life time ~ seconds

Vortex diameter $\sim 60 \ \mu m$

PNAS 2012 & PRL 2013

Cisneros et al, Exp Fluids 2017





Confined bacterial suspensions



Spontaneously broken mirror-symmetry & edge currents

Wioland et al (2013) PRL

Bacterial spin lattices: Anti-ferromagnetic order

Pliit



Wioland et al, Nature Physics 2016

Generalized Navier-Stokes equations

 $\nabla \cdot \boldsymbol{v} = 0,$

 $\partial_t \boldsymbol{v} + \boldsymbol{v} \cdot \nabla \boldsymbol{v} = -\nabla p + \Gamma_0 \nabla^2 \boldsymbol{v} - \Gamma_2 \nabla^4 \boldsymbol{v} + \Gamma_4 \nabla^6 \boldsymbol{v}$



Classical pattern formation theory

Classical turbulence theory

Active liquid crystals



Dogic lab (Brandeis) Nature 2012

Active liquid crystals



Dogic lab, Nature 2012

Active liquid crystal defects



Giomi et al PRL 2012

Topological defects in order-parameter fields

- optical effects
- work hardening, etc





'umbilic defects' in a nematic liquid crystal